

# Independent Measurement of the Top Quark Mass and the Light- and Bottom-Jet Energy Scales at Hadron Colliders

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At hadron-colliders, the measurement of the mass  $m_t$  of the top quark and the absolute energy scale  $S$  for calorimeter jets are closely linked. While the scale  $S_j$  for light jets can be calibrated with hadronic  $W$  decays in the same  $\ell$ +jets  $t\bar{t}$  events used to measure  $m_t$ , the main remaining systematic uncertainty on  $m_t$  is so far due to differences between  $S_j$  and the scale  $S_b$  for  $b$  jets [1]. A novel measurement technique has now been developed that allows a simultaneous determination of  $m_t$ ,  $S_j$ ,  $S_b$ , and the jet energy resolution  $R$  from  $\ell$ +jets  $t\bar{t}$  events [2].

It is assumed that the full calorimeter calibration up to constant scales  $S_j$  and  $S_b$  has been performed before this method is applied. Three *estimators*,  $m_t^{\text{reco}}$ ,  $S_j^{\text{reco}}$ , and  $S_b^{\text{reco}}$ , are calculated for each selected event. Functions are derived to describe the expected estimator distributions (*templates*) for any given set of assumed values of  $m_t$ ,  $S_j$ ,  $S_b$ , and  $R$ . A comparison of the measured estimator distributions in the data with these fitted templates then yields the  $m_t$ ,  $S_j$ ,  $S_b$ , and  $R$  values and their uncertainties.

The method has been tested using  $\ell$ +jets  $t\bar{t}$  events in 14 TeV  $pp$  collisions generated at parton level with ALPGEN [3]. The energies of the final-state quarks have been smeared according to a Gaussian resolution whose width is set to  $\sigma(E) = R\sqrt{E}$  with constant  $R$ . All jet energies are multiplied by a factor  $S_j$ , and  $b$ -jet energies by another factor  $S_b$ . Tests have been performed with various  $(m_t, S_j, S_b, R)$  parameter sets.

Standard  $t\bar{t}$  event selection criteria [2] are applied. In each event, assuming unambiguous  $b$ -jet identification, the estimator  $S_j^{\text{reco}} = \frac{m_W^{\text{raw}}}{m_W}$  is calculated from the known  $W$  mass  $m_W$  and the mass  $m_W^{\text{raw}}$  reconstructed from the smeared light-jet energies. A scan over  $S_b^{\text{reco}}$  values is performed. Given an assumed value of  $S_b^{\text{reco}}$ , the reconstructed  $b$ -quark jet energies and momenta are scaled accordingly, and the missing transverse momentum is adjusted and taken as transverse momentum of the neutrino from the leptonic  $W$  decay. The longitudinal neutrino momentum  $p_\nu^z$  is then obtained from  $m_W$ , and the resulting top quark masses  $m_{t,\text{lep}}^{\text{reco}}$  and  $m_{t,\text{had}}^{\text{reco}}$  of the top quarks with the leptonic/hadronic  $W$  decay are computed. If one finds  $m_{t,\text{lep}}^{\text{reco}} = m_{t,\text{had}}^{\text{reco}}$ , then this top quark mass and the corresponding  $S_b^{\text{reco}}$  value are taken as estimator values for the event. Events are only retained if exactly one solution with  $0.5 < S_b^{\text{reco}} < 2.0$  and  $150 \text{ GeV} < m_t^{\text{reco}} < 200 \text{ GeV}$  is found.

After the preselection, events with a magnitude of the vector sum of  $b$ -quark jet transverse momenta of less than 50 GeV that yield poor independent information on the top quark mass and  $b$ -quark jet energy scale are rejected. Finally, the quantity  $\Delta^{\text{reco}} := \frac{\partial m_{t,\text{lep}}^{\text{reco}}}{\partial S_b^{\text{reco}}} - \frac{\partial m_{t,\text{had}}^{\text{reco}}}{\partial S_b^{\text{reco}}}$  is obtained during the scan of  $S_b^{\text{reco}}$  values. Events with

$\Delta^{\text{reco}} < 30 \text{ GeV}$  have a degraded resolution and are rejected. The resulting  $S_b^{\text{reco}}$  estimator distributions for various choices of input parameters are shown in Figure 1 as an example.

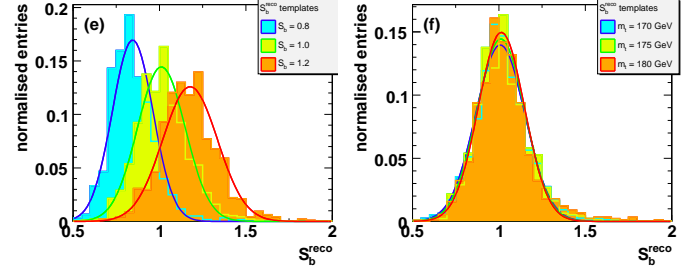


Fig. 1:  $S_b^{\text{reco}}$  template when varying (left plot) the input  $S_b$  value, and (right plot) the input  $m_t$  value. The template parameterizations are overlaid.

To test the method, pseudo-experiments are then performed using simulated events for various sets of input parameter values. Figure 2 shows results for the distributions of measured  $m_t$ ,  $S_j$ , and  $S_b$  values. The correlation matrix between the four measured parameters is given by:

	$S_j$	$S_b$	$R$
$m_t$	-0.09	-0.50	-0.22
$S_j$		-0.38	-0.11
$S_b$			-0.14

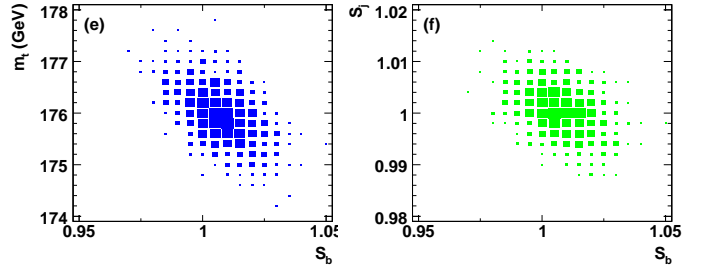


Fig. 2: Pseudo-experiments: The correlation between  $m_t$  and  $S_b$  results (left plot), and that between  $S_j$  and  $S_b$  (right plot).

The parton-level tests of the method have been published [2] as a proof of principle. In the future, the method will be applied to fully simulated ATLAS events, and systematic uncertainties will be investigated.

## References

- [1] F. Fiedler, “Precision Measurements of the Top Quark Mass”, habilitation thesis, 2007, [http://www.etp.physik.uni-muenchen.de/dokumente/thesis/habil\\_ffiedler.pdf](http://www.etp.physik.uni-muenchen.de/dokumente/thesis/habil_ffiedler.pdf)
- [2] F. Fiedler, Eur. Phys. J. C **53** (2008) 41 [arXiv:0706.1640 [hep-ex]].
- [3] M. L. Mangano *et al.*, JHEP **307** (2003) 1